

High energy estimates of the resolvent

Joint work with Thierry Jecko

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Let H be a self-adjoint operator acting on a complex and separable Hilbert space $(\mathcal{H}, \|\cdot\|)$.

We consider *the Schrödinger equation*

$$\begin{cases} i\partial_t f = Hf, \\ f(0) = f_0 \in \mathcal{H}, \end{cases}$$

where $\|f_0\| = 1$.

Let $f(t) = e^{-itH}f_0$ be the unique solution.

If f_0 is an eigenfunction associated to λ_0 , we have : $Hf_0 = \lambda_0 f_0$ and

$$f(t) = e^{-itH}f_0 = e^{-it\lambda_0}f_0.$$

The solution is periodic.

If $f_0 \in \mathcal{H}_{pp}(H)^\perp$, where $\mathcal{H}_{pp}(H) := \overline{\text{span}\{f \text{ eigenfunction of } H\}}$.

Given K which is relatively compact with respect to H , we have :

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \|Ke^{-itH}f_0\|^2 dt = 0.$$

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Aim : If $f_0 \in \mathcal{H}_{pp}(H)^\perp$, where $\mathcal{H}_{pp}(H) := \overline{\text{span}\{f \text{ fonction propre de } H\}}$.

Given K which is relatively compact with respect to H , we have :

$$\lim_{t \rightarrow \infty} Ke^{-itH}f_0 = 0.$$

Example :

In $L^2(\mathbb{R}^n; \mathbb{C})$, set

$$H = -\Delta_{\mathbb{R}^n} + V(Q),$$

with $V \in L^\infty(\mathbb{R}^n; \mathbb{R})$ and where $V(Q)$ is the operator of multiplication by V .

If $\lim_{|x| \rightarrow \infty} V(x) = 0$ and if $f_0 \in \mathcal{H}_{pp}(H)^\perp$, then

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \|\mathbf{1}_{B(0,R)}(Q) e^{-itH} f_0\|^2 dt = 0.$$

If $\lim_{|x| \rightarrow \infty} \langle x \rangle^{1+\varepsilon} V(x) = 0$, for some $\varepsilon > 0$, and if $f_0 \in \mathcal{H}_{pp}(H)^\perp$, then

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We measure the *evolution* with respect to the observable A , i.e. a self-adjoint operator in \mathcal{H} . Set

$$\mathcal{A}_f(t) := \langle e^{-itH}f_0, Ae^{-itH}f_0 \rangle.$$

We want $\mathcal{A}_f(\cdot)$ to be increasing.

We compute the *Heisenberg derivative* :

$$\mathcal{A}'_f(t) = \langle e^{-itH}f_0, [H, iA]e^{-itH}f_0 \rangle.$$

We then suppose that $[H, iA]$ is “strictly positive” (the choice of A is crucial). Then there exists $c > 0$ such that

$$\langle e^{-itH}f_0, Ae^{-itH}f_0 \rangle \geq ct + \langle f_0, Af_0 \rangle$$

There is some *transport*, the particle *diffuses* along A .

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Back to the example : Set

$$A := \frac{1}{2i} (P \cdot Q + Q \cdot P),$$

where $P := -i\nabla$ and Q is the operator of multiplication by x .

We have

$$[-\Delta, iA] = -2\Delta \geq 0.$$

$$E_{\mathcal{I}}(-\Delta)[- \Delta, iA]E_{\mathcal{I}}(-\Delta) = 2(-\Delta)E_{\mathcal{I}}(-\Delta) \geq 2 \inf(\mathcal{I})E_{\mathcal{I}}(-\Delta).$$

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For $\mathcal{I} \in (0, \infty)$.

For the perturbation, we look for some compactness. We suppose

$$V(x) \rightarrow 0 \quad \text{et} \quad x \cdot \nabla V(x) \rightarrow 0,$$

for $|x| \rightarrow \infty$.

This gives

$$\begin{cases} [V, iA] = -Q \cdot \nabla V(Q) \in \mathcal{K}(\mathcal{H}^2, \mathcal{H}), \\ \varphi(-\Delta) - \varphi(H) \in \mathcal{K}(\mathcal{H}, \mathcal{H}^2), \text{ for } \varphi \in C_c^\infty(\mathbb{R}). \end{cases}$$

We deduce that for all interval $\mathcal{I} \in (0, \infty)$, there exist $c > 0$ and K a compact operator, such that

$$E_{\mathcal{I}}(H)[H, iA]E_{\mathcal{I}}(H) \geq \underbrace{c}_{\text{diffusion}} E_{\mathcal{I}}(H) + \underbrace{K}_{\text{e.v.}}$$

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Let \mathcal{I} be a compact interval and A be a self-adjoint operator. We suppose that a *Mourre estimate* holds true (Mourre '81, Perry-Seagal-Simon '81, . . . Amrein-Boutet de Monvel-Georgescu '96) :

$$E_{\mathcal{I}}(H)[H, iA]E_{\mathcal{I}}(H) \geq \underbrace{c}_{\text{const. } > 0} E_{\mathcal{I}}(H) + \underbrace{K}_{\text{compact op.}},$$

in the form sense. We suppose $H \in \mathcal{C}^2(A)$, i.e. that

$$t \mapsto e^{itA}(H+i)^{-1}e^{-itA}$$

is strongly \mathcal{C}^k . Then

- H has at most a finite number of eigenvalues in \mathcal{I} (Virial Theorem).
- Let $\mathcal{J} \in \mathcal{I}$ compact, such that H has no e.v. in \mathcal{J} , then, given K relatively compact with respect to H, we have :

$$\lim_{t \rightarrow \infty} K e^{-itH} E_{\mathcal{J}}(H) f = 0,$$

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- a. H has at most a finite number of eigenvalues in \mathcal{I} (Virial Theorem).
- b'. Let $\mathcal{J} \in \mathcal{I}$ compact, such that H has no eigenvalue in \mathcal{J} , then, there exists C such that

(LAP) $\sup_{\lambda \in \mathcal{J}, \varepsilon > 0} |\langle f, (H - \lambda - i\varepsilon)^{-1} f \rangle| \leq C \|\langle A \rangle^s f\|^2,$
 for all $f \in \mathcal{D}(\langle A \rangle^s)$, $s > 1/2$.

(Kato) $\int_{\mathbb{R}} \|\langle A \rangle^{-s} e^{-itH} E_{\mathcal{J}}(H) f\|^2 dt \leq C \|f\|^2.$

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The Mourre theory has been adapted and used by many authors to deal with thresholds (Fournais-Skibsted, Richard, Bony-Häfner, Bony-Faupin, . . .). We follow the approach of Richard. First we need a different type of positivity.

Back to the example : For $c \in (0, 2)$ and setting $H := -\Delta + V(Q)$, we have

$$[H, iA] - c \times H = (2 - c) \times (-\Delta) - W(Q),$$

where

$$W(Q) := Q \cdot \nabla V(Q) + cV(Q).$$

Assuming also that $W(Q) \geq 0$, we get :

$$[H, iA] - c \times H > 0,$$

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In the abstract setting Richard supposes that there are a self-adjoint operator H and $c \in (0, 2)$ such that

$$[H, iA] - cH > 0.$$

Under some additional regularity assumptions, he infers :

$$\sup_{\lambda \in [0, \infty), \varepsilon > 0} |\langle f, (H - \lambda - i\varepsilon)f \rangle| \leq \|Bf\|^2$$

for some weight B .

Pro :

- ▶ Goes till the threshold.
- ▶ No localisation in energy.
- ▶ The method does not feel the dimension.

Cons :

- ▶ Not the standard regularity assumptions.
- ▶ The weight B is not very natural in some cases.
- ▶ Cannot prove the continuity of the boundary value of the resolvent.
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We adapt the method to study properties of relativistic massive charged particles with spin-1/2. We follow the Dirac formalism. Set :

$$D_m := \alpha \cdot P + m\beta = -i \sum_{k=1}^3 \alpha_k \partial_k + m\beta,$$

where $m > 0$ is the mass. Here we set $\alpha := (\alpha_1, \alpha_2, \alpha_3)$ and $\beta := \alpha_4$. The α_j , for $j \in \{1, 2, 3, 4\}$, are linearly independent self-adjoint linear maps, acting in \mathbb{C}^4 , satisfying the anti-commutation relations :

$$\alpha_i \alpha_j + \alpha_j \alpha_i = 2\delta_{i,j} \mathbf{1}_{\mathbb{C}^4}, \text{ where } i, j \in \{1, 2, 3, 4\}.$$

We add a Coulombic interaction. The Hamiltonian is given by

$$H_\gamma := D_m + \gamma V_C(Q), \text{ where } V_C := v_C \otimes Id_{\mathbb{C}^{2\nu}} \text{ and } v_C(x) := \sum_{k=1, \dots, n} \frac{z_k}{|x - a_k|},$$

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Theorem (Boussaid-G, '10)

There are $\kappa, \delta, C > 0$ such that

$$\sup_{|\lambda| \in [m, m+\delta], \varepsilon > 0, |\gamma| \leq \kappa} \|\langle Q \rangle^{-1} (H_\gamma - \lambda - i\varepsilon)^{-1} \langle Q \rangle^{-1}\| \leq C.$$

In particular, H_γ has no eigenvalue in $\pm m$. Moreover, there is C' so that

$$\sup_{|\gamma| \leq \kappa} \int_{\mathbb{R}} \|\langle Q \rangle^{-1} e^{-itH_\gamma} E_{\mathcal{I}}(H_\gamma) f\|^2 dt \leq C' \|f\|^2,$$

where $\mathcal{I} = [-m - \delta, -m] \cup [m, m + \delta]$ and where $E_{\mathcal{I}}(H_\gamma)$ denotes the spectral measure of H_γ .

Few ideas :

1. Reduce to problem to a bounded perturbation.
2. Use a kind of Feshbach method with respect to the spin-up/down decomposition.
3. Reduce the analyse to a kind of non-selfadjoint Laplacian depending of a parameter.
4. Adapt the method of Richard to this setting (non-self-adjoint + parameter).
5. Conclude by perturbing the resolvent.

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In the present work, we start again with

$$[H, iA] - cH > 0.$$

We want to use the standard class of regularity

and prove that there is $\lambda_0 > 0$ such that

$$\sup_{\lambda \in [\lambda_0, \infty), \varepsilon > 0} |\langle f, (H - \lambda - i\varepsilon)f \rangle| \leq c \|\langle A \rangle f\|^2.$$

Moreover, we would like to cover a non-self-adjoint setting.

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Moreover, we would like to cover a non-self-adjoint setting.

Let H^\pm be two closed and densely defined operators, acting in some Hilbert space \mathcal{H} so that

$$\mathcal{D} := \mathcal{D}(H^+) = \mathcal{D}(H^-) \text{ and } (H^+)^* = H^-.$$

We endow \mathcal{D} with the graph norm of H^+ . In particular, we have that $\mathcal{D}((H^\pm)^*) = \mathcal{D}$.

Since H^\pm are densely defined, share the same domain and are adjoint of the other, we have that $\Re(H^\pm)$ and $\Im(H^\pm)$ are closable operators on \mathcal{D} , indeed their adjoints are densely defined. We denote by $\Re(H^\pm)$ and by $\Im(H^\pm)$ the closure of these operators. It is possible that they are not self-adjoint, albeit there are symmetric. However, \mathcal{D} is a core for them. Their domain is possibly bigger than \mathcal{D} .

We suppose that H^+ is *dissipative*, i.e.

$$\langle f, \Im(H^+)f \rangle \geq 0, \text{ for all } f \in \mathcal{D}.$$

This gives also that $\Im(H^-) \leq 0$. By the numerical range theorem, we infer that $\sigma(H^\pm)$ is included in the half-plan containing $\pm i$.

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Since H^\pm are densely defined, share the same domain and are adjoint of the other, we have that $\Re(H^\pm)$ and $\Im(H^\pm)$ are closable operators on \mathcal{D} , indeed their adjoints are densely defined. We denote by $\Re(H^\pm)$ and by $\Im(H^\pm)$ the closure of these operators. It is possible that they are not self-adjoint, albeit there are symmetric. However, \mathcal{D} is a core for them. Their domain is possibly bigger than \mathcal{D} .

We suppose that H^+ is *dissipative*, i.e.

$$\langle f, \Im(H^+)f \rangle \geq 0, \text{ for all } f \in \mathcal{D}.$$

This gives also that $\Im(H^-) \leq 0$. By the numerical range theorem, we infer that $\sigma(H^\pm)$ is included in the half-plan containing $\pm i$.

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Consider now a non-negative self-adjoint operator S in \mathcal{H} with domain \mathcal{D} . We set $\mathcal{G} := \mathcal{D}(S^{1/2})$. By duality and interpolation we have :

$$H^\pm \in \mathcal{B}(\mathcal{G}, \mathcal{G}^*).$$

Here we identified \mathcal{H} with \mathcal{H}^* .

To perform this analysis, we consider an external operator, the conjugate operator. Let A be a self-adjoint operator in \mathcal{H} . We assume $S \in \mathcal{C}^1(A)$. Let $W_t := e^{itA}$ be the C_0 -group associated to A in \mathcal{H} . We ask :

$$W_t \mathcal{D} \subset \mathcal{D}, \text{ for all } t \in \mathbb{R}.$$

By duality, we have W_t stabilizes \mathcal{D}^* and by interpolation it also stabilizes \mathcal{G} and \mathcal{G}^* .

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Theorem

Let H^\pm and A as above. Suppose that $H^\pm \in \mathcal{C}^2(A; \mathcal{D}, \mathcal{H})$, i.e., there is b s.t.

$$|\langle H^\mp f, Ag \rangle - \langle Af, H^\pm g \rangle| \leq b \|f\| \cdot \|\langle S \rangle g\|, \text{ for all } f, g \in \mathcal{D} \cap \mathcal{D}(A),$$

and

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Assume that, for some $c_1 > 0$,

$$\begin{aligned} [\Re(H^\pm), iA] - c_1 \Re(H^\pm) &\geq S \geq 0, \\ \pm \Im(H^\pm) &\geq 0, \end{aligned}$$

in the sense of forms on \mathcal{G} . Consider $\lambda_0 > 0$. Then, for all $s > 1/2$, there are $c > 0$ and $\mu_0 > 0$, so that

$$|\langle f, (H^\pm - \lambda \pm i\mu)^{-1} f \rangle| \leq c \|\langle |S|^{-1/2} A \rangle^s f\|^2,$$

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Few remarks :

If S is only bounded from below, the theorem is still true if we replace λ_0 big enough.

If $\lambda_0 > 0$ we avoid the threshold effect.

If $c_0 = 0$ one can prove a LAP on \mathbb{R} (see Mantoiu and al).

Note that by duality and interpolation there is $c_2 > 0$ such that

$$|\langle f, [[H^\pm, A], A] f \rangle| \leq c_2 \| (S)^{1/2} f \|^2, \text{ for all } f \in \mathcal{D}.$$

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Example : In $L^2(\mathbb{R}^n; \mathbb{C})$, set

$$H = -\Delta_{\mathbb{R}^n} + V_1(Q) + iV_2(Q),$$

with

$$V_i \text{ and } Q^{1+\varepsilon} \cdot \nabla V_i(Q) \in L^\infty$$

for some $\varepsilon > 0$ and

$$V_2 \geq 0.$$

Given $s > 1/2$, there are $\lambda_0 > 0$ and $c > 0$ such that

$$|\langle f, (H - \lambda - i\mu)^{-1} f \rangle| \leq c \| \langle A \rangle^s f \|^2,$$

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An other example : The bosonic field is described by the Hilbert space $\Gamma(\mathfrak{h})$, where $\mathfrak{h} := L^2(\mathbb{R}^3)$. Set $\mathfrak{h}^{0\otimes} = \mathbb{C}$ and $\mathfrak{h}^{n\otimes} = \mathfrak{h} \otimes \dots \otimes \mathfrak{h}$. Let S_n be the group of permutation of n elements. For each $\sigma \in S_n$, one defines the action on $\mathfrak{h}^{n\otimes}$ by $\sigma(f_{i_1} \otimes \dots \otimes f_{i_n}) = f_{\sigma^{-1}(i_1)} \otimes \dots \otimes f_{\sigma^{-1}(i_n)}$, where (f_i) is a basis of \mathfrak{h} . The action extends to $\mathfrak{h}^{n\otimes}$ by linearity to a unitary operator. The definition is independent of the choice of the basis. On $\mathfrak{h}^{n\otimes}$, we set

$$\Pi_n := \frac{1}{n!} \sum_{\sigma \in S_n} \sigma \text{ and } \Gamma_n(\mathfrak{h}) := \Pi_n(\mathfrak{h}^{n\otimes}). \quad (1)$$

Note that Π_n is an orthogonal projection. We call $\Gamma_n(\mathfrak{h})$ the n -particle bosonic space. The bosonic space is defined by

$$\Gamma(\mathfrak{h}) := \bigoplus_{n=0}^{\infty} \Gamma_n(\mathfrak{h}).$$

We denote by Ω the *vacuum*, the element $(1, 0, 0, \dots)$.

Given a self-adjoint operator ω in \mathfrak{h} and a finite dimensional Hilbert space \mathcal{K} . One defines the free Hamiltonian H_0 acting on the Hilbert space $\mathcal{H} := \mathcal{K} \otimes \Gamma(\mathfrak{h})$ by

$$H_0 := K \otimes 1_{\Gamma(\mathfrak{h})} + 1_{\mathcal{K}} \otimes d\Gamma(\omega), \quad (2)$$

with $\omega := |k|$.

Recall that for b self-adjoint in \mathfrak{h} we define :

$$d\Gamma(b)|_{\Pi_n(\mathcal{D}(b)^{n\otimes})} := \sum_{j=1}^n 1 \otimes \dots \otimes 1 \otimes \underbrace{b}_{j^{th}} \otimes 1 \otimes \dots \otimes 1.$$

We recall also the definition of the *number operator* $N := 1_{\mathcal{K}} \otimes d\Gamma(1)$.

We now define the interaction. Let α be an element $\mathcal{B}(\mathcal{H}, \mathcal{H} \otimes \mathfrak{h})$. This is a *form-factor*. We define $b(\alpha)$ on \mathcal{H} by $b(\alpha) := \mathcal{H} \otimes h^{n \otimes} \rightarrow \mathcal{H} \otimes h^{(n-1) \otimes}$, where

$$b(\alpha)(\Psi \otimes \phi_1 \otimes \dots \otimes \phi_n) := \alpha^*(\Psi \otimes \phi_1) \otimes \phi_2 \otimes \dots \otimes \phi_n,$$

for $n \geq 1$ and by 0 otherwise. This operator is bounded and its norm is given by $\|\alpha\|_{\mathcal{B}(\mathcal{H}, \mathcal{H} \otimes \mathfrak{h})}$. We define the *annihilation operator* on $\mathcal{H} \otimes \Gamma(\mathfrak{h})$ with domain $\mathcal{H} \otimes \Gamma_{fin}(\mathfrak{h})$ by

$$a(\alpha) := (N + 1)^{1/2} b(\alpha)(1 \otimes \Pi),$$

where $\Pi := \sum_n \Pi_n$, see (1). As above, it is closable and its closure is denoted by $a(\alpha)$. Its adjoint is the *creation operator*. It acts as $a^*(\alpha) = b^*(\alpha)(N + 1)^{1/2}$ on \mathcal{H} . Note that $b^*(\alpha)(\psi \otimes \phi_1 \otimes \dots \otimes \phi_n) = (\alpha\psi) \otimes \phi_1 \otimes \dots \otimes \phi_n$. The *Field operator* is defined by

$$\phi(\alpha) := \frac{1}{\sqrt{2}} (a(\alpha) + a^*(\alpha)).$$

We consider its closure on $\mathcal{H} \otimes \mathcal{D}(N^{1/2})$.

Set

$$H_g := H_0 + g\phi(\alpha),$$

and $A := P \cdot Q + Q \cdot P$, where $\omega^{1/2}\alpha$, $\omega^{1/2}A\alpha$, and $\omega^{1/2}A^2\alpha$ are $\mathcal{B}(\mathcal{H}, \mathcal{H} \otimes \mathfrak{h})$. Then given $g_{max} > 0$, there is $\lambda_0 > 0$ we have

$$|\langle f, (H_g - \lambda - \mu)^{-1} f \rangle| \leq \| \langle d\Gamma(A) \rangle^s f \|^2$$

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Partial Proof : We define

$$H_\varepsilon^\pm := H^\pm \pm i\varepsilon[H^\pm, iA],$$

with the common domain \mathcal{D} for $\varepsilon \geq 0$. By Kato-Rellich, since H^\pm is closed and $[H^\pm, iA] \in \mathcal{B}(\mathcal{D}, \mathcal{H})$ imply that, for ε_0 small enough and $|\varepsilon| \leq \varepsilon_0$, H_ε^\pm is closed on the same domain, namely \mathcal{D} .

Since $H^\pm \pm i$ is bijective, by writing

$$H_\varepsilon^\pm \pm i = (1 \pm i\varepsilon[H^\pm, iA] \circ (H^\pm \pm i)^{-1})(H^\pm \pm i),$$

as operators acting on \mathcal{D} , and using again $[H^\pm, iA] \in \mathcal{B}(\mathcal{D}, \mathcal{H})$, we see that $H_\varepsilon^\pm \pm i$ is bijective (and therefore closed) for ε small enough. Hence $(H_\varepsilon^\pm \pm i)^*$ is also bijective from $\mathcal{D}((H_\varepsilon^\pm)^*)$ onto \mathcal{H} .

Now since $(H_\varepsilon^\pm \pm i)^*$ is a bijective extension of the bijective map $H_\varepsilon^\mp \mp i$ (from \mathcal{D} onto \mathcal{H}), $(H_\varepsilon^\pm \pm i)^* = H_\varepsilon^\mp \mp i$. Therefore, for $|\varepsilon| \leq \varepsilon_0$, $\mathcal{D}((H_\varepsilon^\pm)^*) = \mathcal{D}$ and $(H_\varepsilon^\pm)^* = H_\varepsilon^\mp$.

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Since $H^\pm \in C^1(A; \mathcal{G}, \mathcal{G}^*)$, we obtain that

$$[H^\pm, A] = [\Re(H^\pm), A] + i[\Im(H^\pm), A] \in \mathcal{B}(\mathcal{G}, \mathcal{G}^*)$$

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$$H_\varepsilon^\pm - \lambda \pm i\mu \in \mathcal{B}(\mathcal{G}, \mathcal{G}^*).$$

Now, take $f \in \mathcal{G}$. Take $\varepsilon, \lambda, \mu \geq 0$. We get :

$$\begin{aligned} & -c_1\varepsilon \langle f, \Re(H_\varepsilon^\pm - \lambda \pm i\mu)f \rangle \pm \langle f, \Im(H_\varepsilon^\pm - \lambda \pm i\mu)f \rangle = \\ & = -c_1\varepsilon \langle f, (\Re(H^\pm) \mp \varepsilon[\Im(H^\pm), iA] - \lambda) f \rangle \pm \langle f, (\Im(H^\pm) \pm \mu \pm \varepsilon[\Re(H^\pm), iA]) f \rangle \\ & = \varepsilon \langle f, ([\Re(H^\pm), iA] - c_1\Re(H^\pm))f \rangle + (c_1\lambda\varepsilon + \mu) \|f\|^2 \\ & \quad \pm \langle f, (c_1\varepsilon^2[\Im(H^\pm), iA] + \Im(H^\pm)) f \rangle \\ & \geq (c_1\lambda\varepsilon + \mu) \|f\|^2 + \varepsilon \|S^{1/2}f\|^2. \end{aligned} \tag{3}$$

Here we remove the condition on the commutator of the imaginary part by taking ε small enough.

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Here we remove the condition on the commutator of the imaginary part by taking ε small enough.

We start with a crude bound. For $f \in \mathcal{G}$, we set :

$$\|f\|_{\mathcal{G}}^2 := \|f\|^2 + \|\mathcal{S}^{1/2}f\|^2.$$

Using Cauchy-Schwartz, we obtain :

$$\sqrt{2} \max(c_1 \varepsilon, 1) \|(H_\varepsilon^\pm - \lambda \pm i\mu)f\|_{\mathcal{G}^*} \geq \min(c_1 \lambda \varepsilon + \mu, \varepsilon) \|f\|_{\mathcal{G}}.$$

This shows that, as an operator in $\mathcal{B}(\mathcal{G}, \mathcal{G}^*)$, $H_\varepsilon^\pm - \lambda \pm i\mu$ is injective with closed range. As such, its adjoint belongs to $\mathcal{B}(\mathcal{G}, \mathcal{G}^*)$ and is an extension of $H_\varepsilon^\mp - \lambda \mp i\mu$. Thus it is actually $H_\varepsilon^\mp - \lambda \mp i\mu$, which is injective with closed range. All this implies the bijectivity of $H_\varepsilon^\pm - \lambda \pm i\mu$, viewed in $\mathcal{B}(\mathcal{G}, \mathcal{G}^*)$, and

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Then, we first restrict the domain of G_ε^\pm to \mathcal{H} and improve it.

Since (3) holds true on \mathcal{D} , the common domain of H_ε^\pm and $(H_\varepsilon^\pm)^*$, we can apply the numerical range theorem. Thus the spectrum of $H_\varepsilon^\pm - \lambda \pm i\mu$ is contained in the half-plane delimited by the equation

$$\pm y \geq c_1 \varepsilon x + \mu.$$

Taking possibly ε_0 smaller (but positive), the distance from 0 to this half-plane is larger than $\mu/2$. Then $G_\varepsilon^\pm|_{\mathcal{H}} = (H_\varepsilon^\pm - \lambda \pm i\mu)^{-1}$ exists from \mathcal{H} to \mathcal{D} . Thus

$$\|G_\varepsilon^\pm\|_{B(\mathcal{H})} \leq 2/\mu, \text{ for } \mu > 0 \text{ and } \varepsilon \in [0, \varepsilon_0]. \quad (5)$$

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$$F_\varepsilon^\pm := \langle f, G_\varepsilon^\pm f \rangle.$$

Notice that $\overline{F_\varepsilon^\pm} = F_\varepsilon^\mp$. Since $G_\varepsilon^\pm \mathcal{H} = \mathcal{D} \subset \mathcal{G}$ and using (4), we infer

$$\min(c_1 \lambda \varepsilon + \mu, \varepsilon) \left\| \langle S \rangle^{1/2} G_\varepsilon^\pm f \right\|^2 \leq \sqrt{2} \max(c_1 \varepsilon, 1) |F_\varepsilon^\pm|. \quad (6)$$

Hence up to a smaller $\varepsilon_0 > 0$ and since $\mu > 0$, we obtain

$$\left\| \langle S \rangle^{1/2} G_\varepsilon^\pm f \right\|^2 \leq \frac{1}{\varepsilon \min(c_1 \lambda, 1)} |F_\varepsilon^\pm| = \frac{c_0}{\varepsilon} |F_\varepsilon^\pm|, \text{ for all } \varepsilon \in (0, \varepsilon_0], \quad (7)$$

and where $c_0 := 1/\min(1, \lambda_0 c_1) \geq 1$. From (6), we derive

$$|F_\varepsilon^\pm| \leq \left\| \langle S \rangle^{-1/2} f \right\| \left\| \langle S \rangle^{1/2} G_\varepsilon^\pm f \right\| \leq \sqrt{\frac{c_0}{\varepsilon}} \left\| \langle S \rangle^{-1/2} f \right\| \sqrt{|F_\varepsilon^\pm|}$$

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We now show that $G_\varepsilon^\pm \in \mathcal{C}^1(A)$. Recall that G_ε^\pm is a bijection from \mathcal{H} onto \mathcal{D} . Then by taking the adjoint, it is also a bijection from \mathcal{D}^* onto \mathcal{H} . Remember now that e^{itA} stabilizes \mathcal{G} and \mathcal{G}^* . By the resolvent equality in $\mathcal{B}(\mathcal{H})$, we have :

$$\frac{1}{t}[G_\varepsilon^\pm, e^{itA}] = -\frac{1}{t} \underbrace{G_\varepsilon^\pm}_{\mathcal{H} \leftarrow \mathcal{G}^*} \underbrace{[H^\pm \pm i\varepsilon[H^\pm, iA], e^{itA}]}_{\mathcal{G}^* \leftarrow \mathcal{G}} \underbrace{G_\varepsilon^\pm}_{\mathcal{G} \leftarrow \mathcal{H}}$$

Let us now take the limit at 0. Since H^\pm and $[H^\pm, iA]$ are in $\mathcal{C}^1(A; \mathcal{G}, \mathcal{G}^*)$, the r.h.s has a strong limit for all element in \mathcal{H} . Hence, $G_\varepsilon^\pm \in \mathcal{C}^1(A; \mathcal{H}, \mathcal{H})$ which actually means $G_\varepsilon^\pm \in \mathcal{C}^1(A)$.

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Considering commutators in the form sense, we have :

$$\frac{dF_\varepsilon^\pm}{d\varepsilon} = \left\langle f, \frac{dG_\varepsilon^\pm}{d\varepsilon} f \right\rangle \quad (9)$$

$$\begin{aligned} &= \pm i \langle G_\varepsilon^\mp f, [H^\pm, iA] G_\varepsilon^\pm f \rangle \\ &= \pm i \langle G_\varepsilon^\mp f, [H_\varepsilon^\pm - \lambda \pm i\mu, iA] G_\varepsilon^\pm f \rangle - \varepsilon \langle G_\varepsilon^\mp f, [[H^\pm, iA], iA] G_\varepsilon^\pm f \rangle \\ &= \pm \langle G_\varepsilon^\mp f, Af \rangle \mp \langle Af, G_\varepsilon^\pm f \rangle - \varepsilon \langle G_\varepsilon^\mp f, [[H^\pm, iA], iA] G_\varepsilon^\pm f \rangle, \end{aligned} \quad (10)$$

for all $f \in \mathcal{D}(A)$.

Using the hypothesis on the second commutator, we infer :

$$\begin{aligned} \left\| \frac{dF_\varepsilon^\pm}{d\varepsilon} \right\| &\leq \left\| \langle S \rangle^{-1/2} A f \right\| \left(\left\| \langle S \rangle^{1/2} G_\varepsilon^\pm f \right\| + \left\| \langle S \rangle^{1/2} G_\varepsilon^\mp f \right\| \right) \\ &\quad + c_2 \varepsilon \left\| \langle S \rangle^{1/2} G_\varepsilon^\pm f \right\| \cdot \left\| \langle S \rangle^{1/2} G_\varepsilon^\mp f \right\|. \end{aligned} \quad (11)$$

Now we use four times (7) and the fact that $|F_\varepsilon^\pm| = |F_\varepsilon^\mp|$, and integrate to obtain, for all $f \in \mathcal{D}(A)$,

$$\left| F_{\varepsilon'}^\pm - F_\varepsilon^\pm \right| \leq \int_\varepsilon^{\varepsilon'} \left\{ 2s^{-1/2} \sqrt{c_0} \sqrt{|F_s^\pm|} \left\| \langle S \rangle^{-1/2} A f \right\| + c_2 c_0 |F_s^\pm| \right\} ds, \quad (12)$$

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By the Gronwall lemma, we infer that there is $C_1 > 0$, independent of $c_0 \geq 1$, $\varepsilon \in (0, \varepsilon_0]$, $\lambda \geq \lambda_0$, and of $\mu > 0$, such that, for $f \in \mathcal{D}(A)$,

$$\begin{aligned}
 |F_\varepsilon^\pm| &\leq e^{c_0 c_2 (\varepsilon_0 - \varepsilon)} \left(|F_{\varepsilon_0}^\pm|^{1/2} + \int_\varepsilon^{\varepsilon_0} \left\{ \frac{1}{\sqrt{\eta}} e^{-\frac{1}{2} c_0 c_2 (\varepsilon_0 - \eta)} \right\} d\eta \left\| \langle S \rangle^{-1/2} A f \right\| \sqrt{c_0} \right)^2 \\
 &\leq 2e^{c_0 c_2 \varepsilon_0} \left(|F_{\varepsilon_0}^\pm| + C_1 c_0 \left\| \langle S \rangle^{-1/2} A f \right\|^2 \right) \\
 &\leq 2c_0 e^{c_0 c_2 \varepsilon_0} \left(\varepsilon_0^{-1} \left\| \langle S \rangle^{-1/2} f \right\|^2 + C_1 \left\| \langle S \rangle^{-1/2} A f \right\|^2 \right). \tag{13}
 \end{aligned}$$

using (8).

We now plug this back into (12). Thus there exist $C_2, C_3 > 0$ with the same independence such that

$$\begin{aligned} \left| F_{\varepsilon'}^{\pm} - F_{\varepsilon}^{\pm} \right| &\leq c_0^2 e^{c_0 c_2 \varepsilon_0} \int_{\varepsilon}^{\varepsilon'} \left\{ 2 \frac{\sqrt{C_2}}{\sqrt{s}} + C_1 C_2 \right\} ds (\|f\|^2 + \|\langle S \rangle^{-1/2} A f\|^2) \\ &= C_3 c_0^2 e^{c_0 c_2 \varepsilon_0} (\sqrt{\varepsilon'} - \sqrt{\varepsilon}) (\|f\|^2 + \|\langle S \rangle^{-1/2} A f\|^2). \end{aligned}$$

Then, $\{F_{\varepsilon}^{\pm}\}_{\varepsilon \in (0, \varepsilon_0]}$ is a Cauchy sequence. We denote by F_{0+}^{\pm} the limit, as ε goes to 0. It remains to notice that $F_{0+}^{\pm} = F_0^{\pm}$. Indeed, using the resolvent identity, (5), and (1),

$$\|G_0^{\pm} - G_{\varepsilon}^{\pm}\|_{\mathcal{B}(\mathcal{H})} \leq \varepsilon \|G_{\varepsilon}^{\pm}\|_{\mathcal{B}(\mathcal{H})} \cdot \|[H^{\pm}, iA](H^{\pm} - \lambda \pm i\mu)^{-1}\|_{\mathcal{B}(\mathcal{H})} \leq \frac{c(\lambda)\varepsilon}{\mu^2},$$

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